

# Assignment of the Side-Chain $^1\text{H}$ and $^{13}\text{C}$ Resonances of Interleukin-1 $\beta$ Using Double- and Triple-Resonance Heteronuclear Three-Dimensional NMR Spectroscopy<sup>†</sup>

G. Marius Clore,\*<sup>†</sup> Ad Bax,<sup>†</sup> Paul C. Driscoll,<sup>†</sup> Paul T. Wingfield,<sup>§</sup> and Angela M. Gronenborn\*<sup>†</sup>

Laboratory of Chemical Physics, Building 2, National Institute of Diabetes and Digestive and Kidney Diseases, and Protein Expression Laboratory, Building 6B, National Institutes of Health, Bethesda, Maryland 20892

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**ABSTRACT:** The assignment of the aliphatic  $^1\text{H}$  and  $^{13}\text{C}$  resonances of IL-1 $\beta$ , a protein of 153 residues and molecular mass 17.4 kDa, is presented by use of a number of novel three-dimensional (3D) heteronuclear NMR experiments which rely on large heteronuclear one-bond  $J$  couplings to transfer magnetization and establish through-bond connectivities. These 3D NMR experiments circumvent problems traditionally associated with the application of conventional 2D  $^1\text{H}$ - $^1\text{H}$  correlation experiments to proteins of this size, in particular the extensive chemical shift overlap which precludes the interpretation of the spectra and the reduced sensitivity arising from  $^1\text{H}$  line widths that are often significantly larger than the  $^1\text{H}$ - $^1\text{H}$   $J$  couplings. The assignment proceeds in two stages. In the first step the  $^{13}\text{C}\alpha$  chemical shifts are correlated with the NH and  $^{15}\text{N}$  chemical shifts by a 3D triple-resonance NH- $^{15}\text{N}$ - $^{13}\text{C}\alpha$  (HNCA) correlation experiment which reveals both intraresidue NH( $i$ )- $^{15}\text{N}(i)$ - $^{13}\text{C}\alpha(i)$  and some weaker interresidue NH( $i$ )- $^{15}\text{N}(i)$ -C $\alpha(i-1)$  correlations, the former via intraresidue one-bond  $^1J_{\text{NC}\alpha}$  and the latter via interresidue two-bond  $^2J_{\text{NC}\alpha}$  couplings. As the NH,  $^{15}\text{N}$ , and C $\alpha$ H chemical shifts had previously been sequentially assigned by 3D  $^1\text{H}$  Hartmann-Hahn  $^{15}\text{N}$ - $^1\text{H}$  multiple quantum coherence (3D HOHAHA-HMQC) and 3D heteronuclear  $^1\text{H}$  nuclear Overhauser  $^{15}\text{N}$ - $^1\text{H}$  multiple quantum coherence (3D NOESY-HMQC) spectroscopy [Driscoll, P. C., Clore, G. M., Marion, D., Wingfield, P. T., & Gronenborn, A. M. (1990) *Biochemistry* 29, 3542-3556], the 3D triple-resonance HNCA correlation experiment permits the sequence-specific assignments of  $^{13}\text{C}\alpha$  chemical shifts in a straightforward manner. The second step involves the identification of side-chain spin systems by 3D  $^1\text{H}$ - $^{13}\text{C}$ - $^{13}\text{C}$ - $^1\text{H}$  correlated (HCCH-COSY) and 3D  $^1\text{H}$ - $^{13}\text{C}$ - $^{13}\text{C}$ - $^1\text{H}$  total correlated (HCCH-TOCSY) spectroscopy, the latter making use of isotropic mixing of  $^{13}\text{C}$  magnetization to obtain relayed connectivities along the side chains. Extensive cross-checks are provided in the assignment procedure by examination of the connectivities between  $^1\text{H}$  resonances at all the corresponding  $^{13}\text{C}$  shifts of the directly bonded  $^{13}\text{C}$  nuclei. In this manner, we were able to obtain complete  $^1\text{H}$  and  $^{13}\text{C}$  side-chain assignments for all residues, with the exception of 4 (out of a total of 15) lysine residues for which partial assignments were obtained. The 3D heteronuclear correlation experiments described are highly sensitive, and the required set of three 3D spectra was recorded in only 1 week of measurement time on a single uniformly  $^{15}\text{N}/^{13}\text{C}$ -labeled 1.7 mM sample of interleukin-1 $\beta$ . This first example of the essentially complete side-chain assignments of a protein with a molecular mass greater than 15 kDa by heteronuclear 3D NMR methods provides a basis for the determination of a full high-resolution three-dimensional structure of interleukin-1 $\beta$  in solution.

Interleukin-1 $\beta$  (IL-1 $\beta$ )<sup>1</sup> is a member of the cytokine family of proteins, which play a central role in the immune and inflammatory responses. It is a protein of 153 amino acids with a molecular mass of 17.4 kDa, and its wide-ranging biological activities include stimulation of B-lymphocyte proliferation and fever induction [see Oppenheim et al. (1986), Dinarello (1988), and Moore (1989) for reviews]. In two recent papers we presented the complete assignment of the  $^{15}\text{N}$  and  $^1\text{H}$  resonances of the polypeptide backbone of IL-1 $\beta$  and the determination of its secondary structure and molecular topology using 3D heteronuclear  $^1\text{H}$ - $^{15}\text{N}$  NMR spectroscopy (Driscoll et al., 1990a,b). At that time only a limited number of  $^1\text{H}$  side-chain resonance assignments were obtained owing to two factors. First, the sensitivity of homonuclear  $J$  correlation experiments (COSY, P.COSY, DQF-COSY, HOHAHA, etc.) is reduced in larger proteins owing to their slow

tumbling, which results in  $^1\text{H}$  line widths that are often significantly larger than many of the  $^1\text{H}$ - $^1\text{H}$  couplings. As a result, only a limited number of relayed through-bond connectivities were observed in the  $^1\text{H}$ - $^{15}\text{N}$  3D HOHAHA-HMQC experiments, thereby preventing in the majority of cases the correlation of NH chemical shifts with those of side-chain protons (C $\beta$ H and beyond). Second, extensive chemical shift overlap in the aliphatic region of the 2D  $^1\text{H}$ - $^1\text{H}$

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<sup>†</sup> Laboratory of Chemical Physics.

<sup>§</sup> Protein Expression Laboratory.

<sup>1</sup> Abbreviations: IL-1 $\beta$ , interleukin-1 $\beta$ ; COSY, correlated spectroscopy; HOHAHA, homonuclear Hartmann-Hahn spectroscopy; TOCSY, total correlated spectroscopy; HMQC, heteronuclear multiple quantum correlation spectroscopy; DQF-COSY, double quantum filtered correlated spectroscopy; P.COSY, primitive correlated spectroscopy; HNCA, 3D triple-resonance NH- $^{15}\text{N}$ - $^{13}\text{C}\alpha$  correlation spectroscopy; HCCH-COSY, 3D  $^1\text{H}$ - $^{13}\text{C}$ - $^{13}\text{C}$ - $^1\text{H}$  correlation spectroscopy via  $^1J_{\text{CC}}$  carbon couplings; HCCH-TOCSY, 3D  $^1\text{H}$ - $^{13}\text{C}$ - $^{13}\text{C}$ - $^1\text{H}$  total correlation spectroscopy with isotropic mixing of  $^{13}\text{C}$  magnetization;  $^1\text{H}$ - $^{13}\text{C}$  NOESY-HMQC, 3D heteronuclear  $^1\text{H}$  nuclear Overhauser  $^{13}\text{C}$ - $^1\text{H}$  multiple quantum coherence spectroscopy;  $^1\text{H}$ - $^{15}\text{N}$  NOESY-HMQC, 3D heteronuclear  $^1\text{H}$  nuclear Overhauser  $^{15}\text{N}$ - $^1\text{H}$  multiple quantum coherence spectroscopy;  $^1\text{H}$ - $^{15}\text{N}$  HOHAHA-HMQC, 3D heteronuclear  $^1\text{H}$  Hartmann-Hahn  $^{15}\text{N}$ - $^1\text{H}$  multiple quantum coherence spectroscopy; NOE, nuclear Overhauser effect; rf, radiofrequency.

## EXPERIMENTAL PROCEDURES

**NMR Spectroscopy.** All NMR experiments were carried out at 36 °C on either a Bruker AM600 or a Bruker AM500 spectrometer, operating in the “reverse” mode. The pulse sequences used for the 3D triple-resonance HN-<sup>15</sup>N-<sup>13</sup>Cα correlation experiment (HNCA; Ikura et al., 1990a; Kay et al., 1990b) and the 3D double-resonance HCCH-COSY (Kay et al., 1990a; Bax et al., 1990a) and HCCH-TOCSY (Bax et al., 1990b) experiments have been described in detail previously and therefore will only be summarized below.

The pulse sequence for the 3D HCCH-COSY experiment is

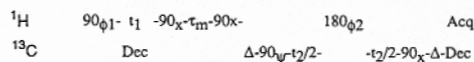
$$\begin{array}{l} {}^1\text{H} \quad 180_x \quad -\text{Acq} \\ {}^{13}\text{C} \quad -\tau-180_{\phi\phi}-\tau-90_x-90_{\phi\phi}-\text{Dec} \end{array}$$

$^1\text{H}$	180 $^\circ$	-Acc
$^{13}\text{C}$	- $\tau$ -180 $^\circ$ - $\tau$ -90 $^\circ$ -90 $^\circ$ -Dec	

The phase cycling for both experiment is as follows:  $\phi_1 = y, -y$ ;  $\phi_2 = 4(x), 4(y), 4(-x), 4(-y)$ ;  $\phi_3 = 8(x), 8(-x)$ ;  $\phi_4 = 2(x), 2(-x)$ ;  $\phi_5 = 2(x), 2(y), 2(-x), 2(-y)$ ;  $\phi_6 = 4(x), 4(-x)$ ;  $\phi_7 = 8(x), 8(y)$ ;  $\text{Acq} = 2(x, -x, -x, x), 2(-x, x, x, -x)$ . Quadrature detection in the  $t_1$  and  $t_2$  dimensions is obtained with the TPPI-States method (Marion et al., 1989) using  $\psi_1 = 16(x), 16(y)$  and  $\psi_2 = 32(x), 32(y)$ . Each time  $t_1$  is incremented, the receiver reference phase and  $\psi_1$  are incremented by  $180^\circ$ ; similarly, each time  $t_2$  is incremented, the receiver phase and  $\psi_2$  are incremented by  $180^\circ$ . Data obtained for  $\psi_1 = x, y$  and  $\psi_2 =$

$x, y$  are stored separately and processed as complex data. The  $180_{\phi 3}$  and  $180_{\phi 4}$  pulses are composite  $180^\circ$  pulses (of the  $90_x 180_y 90_x$  variety). The  $^1\text{H}$  and  $^{13}\text{C}$  carriers were positioned at 3 and 43 ppm, respectively. The carbonyl  $180_{\phi 7}$  pulse used to decouple the carbonyls is applied as a DANTE series (Kay et al., 1989) consisting of eight cycles of six  $4^\circ$ – $5^\circ$  pulses, with the phase of each successive pulse decremented by  $60^\circ$ . The duration of each cycle is set to  $1/\delta$  ( $\sim 50$ – $60$   $\mu\text{s}$ ), where  $\delta$  is the offset of the carbonyl resonances from the  $^{13}\text{C}$  rf carrier. In the case of the HCCH-TOCSY experiment, the duration of the spin-lock trim pulse (SL) is 2 ms, and all the pulses of the DIPSI-3 isotropic mixing sequence (Shaka et al., 1988) are applied along the  $\pm y$  axis with a 7-kHz rf field strength. The length used for the DIPSI-3 isotropic mixing period was 24 ms. The delay  $\tau$  is set to 1.5 ms, slightly less than  $1/(4J_{\text{CH}})$ ; the delays  $\delta 1$  and  $\delta 2$  are set to  $\sim 1/(6J_{\text{CH}}) \sim 1.1$  ms, and the delay  $\Delta + \delta 1$  is set to  $\sim 1/(8J_{\text{CC}}) \sim 3.25$  ms. To eliminate the need for phase correction in the  $F_1$  dimension, the duration of the composite  $180^\circ$   $^{13}\text{C}$  pulse (100  $\mu\text{s}$ ) was added to the  $\tau$  delay between the  $180_x$  and  $90_{\phi 1}$   $^1\text{H}$  pulses. In the  $F_2$  dimension, phase correction is eliminated by applying the  $180_{\phi 2}$   $^1\text{H}$  pulse at the end of the first  $\delta 1$  interval.  $^{13}\text{C}$  decoupling during the acquisition period ( $t_3$ ) is achieved with GARP (Shaka et al., 1985) modulation. The HCCH-COSY experiment was recorded at 600 MHz while the HCCH-TOCSY one was recorded at 500 MHz. A total of 128 complex  $t_1$  ( $^1\text{H}$ ), 32 complex  $t_2$  ( $^{13}\text{C}$ ), and 512 real  $t_3$  ( $^1\text{H}$ ) data points were collected for both experiments with acquisition times of 28.16 ( $t_1$ ), 10.24 ( $t_2$ ) and 47.62 ( $t_3$ ) ms for the HCCH-COSY experiment at 600 MHz and acquisition times of 33.79 ( $t_1$ ), 12.29 ( $t_2$ ), and 51.2 ( $t_3$ ) ms for the HCCH-TOCSY experiment at 500 MHz. Thus, the spectral widths in the  $F_1$  and  $F_2$  dimensions were the same for both experiments, 7.57 and 20.71 ppm, respectively, while the  $F_3$  spectral width was 8.96 ppm for the HCCH-COSY experiment and 10 ppm for the HCCH-TOCSY experiment. The total measurement time for each experiment was approximately 70 h. Because the  $F_2$  spectral width for both experiments was set to 20.71 ppm, extensive folding occurs in this dimension, and each  $^{13}\text{C}$  slice contains resonances corresponding to three different  $^{13}\text{C}$  chemical shifts.

In addition to the three correlation experiments, we also recorded a  $^1\text{H}$ – $^{13}\text{C}$  NOESY-HMQC experiment (Ikura et al., 1990b; Zuiderweg et al., 1990) at 600 MHz with a 130-ms ( $\tau_m$ ) NOESY mixing time using the pulse sequence



with phases  $\phi 2 = 4(x), 4(y)$  and  $\text{Acq} = x, 2(-x), x, -x, 2(x), -x$ . Quadrature in  $F_1$  and  $F_2$  is obtained with the TPPI-States method (Marion et al., 1989), with  $\phi 1 = 4(x, -x), 4(y, -y)$  and  $\psi = 4[2(x), 2(-x)], 4[2(y), 2(-y)]$  and data for  $\phi 1 = x, y$  and  $\psi = x, y$  being stored separately to be processed as complex data. Each time  $t_1$  is incremented,  $\phi 1$  and the receiver phase are incremented by  $180^\circ$ ; and similarly, each time  $t_2$  is incremented,  $\psi$  and the receiver phase are incremented by  $180^\circ$ .  $^{13}\text{C}$  decoupling during the  $t_1$  and acquisition periods is achieved with GARP (Shaka et al., 1985) modulation. The delay  $\Delta$  is set to 3 ms, slightly less than  $1/(2J_{\text{CH}})$ . For convenience the initial  $t_2$  delay is adjusted such that the phase correction in  $F_2$  is  $-180^\circ$  for the first-order phase and  $90^\circ$  for the zero-order phase by setting  $\Delta t_2/2 = [\tau_{180}(^1\text{H}) + 4\tau_{90}(^{13}\text{C})]/\pi + t_2(0)$  by appropriate adjustment of  $t_2(0)$  (where  $\Delta t_2$  is the increment used in the  $t_2$  dimension). The  $^1\text{H}$  carrier was placed at 3.5 ppm and the  $^{13}\text{C}$  carrier at 71.8 ppm. A total of 128 complex  $t_1$  ( $^1\text{H}$ ), 64 complex  $t_2$  ( $^{13}\text{C}$ ), and 512 real  $t_3$  ( $^1\text{H}$ ) data points were collected with acquisition times of 24.32 ( $t_1$ ), 10.24

( $t_2$ ), and 47.62 ( $t_3$ ) ms, corresponding to sweep widths of 8.76 ( $F_1$ ), 41.42 ( $F_2$ ), and 10.41 ( $F_3$ ) ppm. Thus, the sweep width employed in the  $^{13}\text{C}$   $F_2$  dimension was exactly double that used for the HCCH-COSY and HCCH-TOCSY experiments. By this means, any possible ambiguities in  $^{13}\text{C}$  chemical shifts arising from folding are easily resolved by comparing the two sets of experiments. In addition, the NOESY-HMQC spectrum provides an easy means of assigning the  $^{13}\text{C}$  chemical shifts of aromatic carbons bound to protons.

**Processing of 3D Spectra.** All 3D NMR experiments were processed on a Sun Sparc workstation using a simple in-house-written routine for the  $F_2$  Fourier transform together with the commercially available software package NMR2 (New Methods Research, Inc., Syracuse, NY) for processing the  $F_1$ – $F_3$  planes, as described previously (Kay et al., 1989). Zero-filling (once in each dimension) was employed to yield a final absorptive spectrum of  $128 \times 64 \times 512$  data points for the HNCA spectrum,  $256 \times 64 \times 512$  data points for the HCCH-COSY and HCCH-TOCSY spectra, and  $256 \times 128 \times 512$  data points for the NOESY-HMQC spectrum. The weighting function applied in the  $F_1$  and  $F_3$  dimensions was a singly  $60^\circ$ -shifted sine bell window function, while a doubly shifted sine bell function, shifted by  $60^\circ$  at the beginning of the window and  $10^\circ$  at the end of the window, was used in the  $F_2$  dimension (Kay et al., 1989; Driscoll et al., 1990a).

## RESULTS AND DISCUSSION

The NH,  $\text{C}^\alpha\text{H}$ , and  $^{15}\text{N}$  chemical shifts had previously been assigned unambiguously by the 3D  $^1\text{H}$ – $^{15}\text{N}$  HOHAHA-HMQC and NOESY-HMQC experiments (Driscoll et al., 1990a). Because of extensive overlap of  $\text{C}^\alpha\text{H}$  chemical shifts, a helpful first step in the assignment of side chains involves the sequence-specific identification of  $^{13}\text{C}\alpha$  chemical shifts. For larger proteins with many degenerate  $\text{C}^\alpha\text{H}$  chemical shifts, the  $^{13}\text{C}\alpha$  shifts cannot simply be determined by recording a 2D  $^1\text{H}$ – $^{13}\text{C}$  shift correlation spectrum (Brühweiler & Wagner, 1986), and a more sophisticated approach is required. To correlate the  $\text{C}^\alpha\text{H}$  and  $^{13}\text{C}\alpha$  chemical shifts of a given residue unambiguously, we proceeded to use the 3D NH– $^{15}\text{N}$ – $^{13}\text{C}\alpha$  (HNCA) triple-resonance correlation experiment, which links the intrareidue NH,  $^{15}\text{N}$ , and  $^{13}\text{C}\alpha$  chemical shifts via the large (95 Hz) one-bond  $^1J_{\text{NH}}$  coupling and the relatively small (8–12 Hz) one-bond  $^1J_{\text{NC}\alpha}$  intrareidue coupling (Ikura et al., 1990a; Kay et al., 1990b). Briefly, the HNCA experiment transfers magnetization originating on an NH proton to its directly bonded  $^{15}\text{N}$  spin via an INEPT sequence, followed by the evolution of  $^{15}\text{N}$  chemical shifts during the period  $t_1$ . Subsequent application of  $90^\circ$  pulses to both  $^1\text{H}$  and  $^{13}\text{C}\alpha$  spins establishes three-spin NH– $^{15}\text{N}$ – $^{13}\text{C}\alpha$  coherence, and the evolution of solely  $^{13}\text{C}\alpha$  chemical shifts during the period  $t_2$  is ensured by refocusing of  $^1\text{H}$  and  $^{15}\text{N}$  chemical shifts through the application of  $^1\text{H}$  and  $^{15}\text{N}$   $180^\circ$  pulses at the midpoint of the  $t_2$  period. Magnetization is then transferred back to the NH protons by simply reversing the above procedure. The end result is a 3D spectrum in which each peak is labeled by  $^{15}\text{N}$ ,  $^{13}\text{C}\alpha$ , and NH chemical shifts in the  $F_1$ ,  $F_2$ , and  $F_3$  dimensions, respectively.

Several slices of the 3D HNCA spectrum at different  $^{15}\text{N}$ –( $F_1$ ) chemical shifts are shown in Figure 1. It will be noted that, in addition to the intrareidue NH( $i$ )– $^{15}\text{N}(i)$ – $\text{C}\alpha(i)$  correlations, there are also weaker interresidue NH( $i$ )– $^{15}\text{N}(i-1)$ – $\text{C}\alpha(i-1)$  correlations which arise via the small ( $\leq 7$  Hz) two-bond  $^2J_{\text{NC}\alpha}$  interresidue coupling. A total of 61 such sequential connectivities were observed, all of which confirmed the previous assignments. Further, these interresidue two-bond  $J$  connectivities enabled us to obtain the  $^{13}\text{C}\alpha$  chemical shift



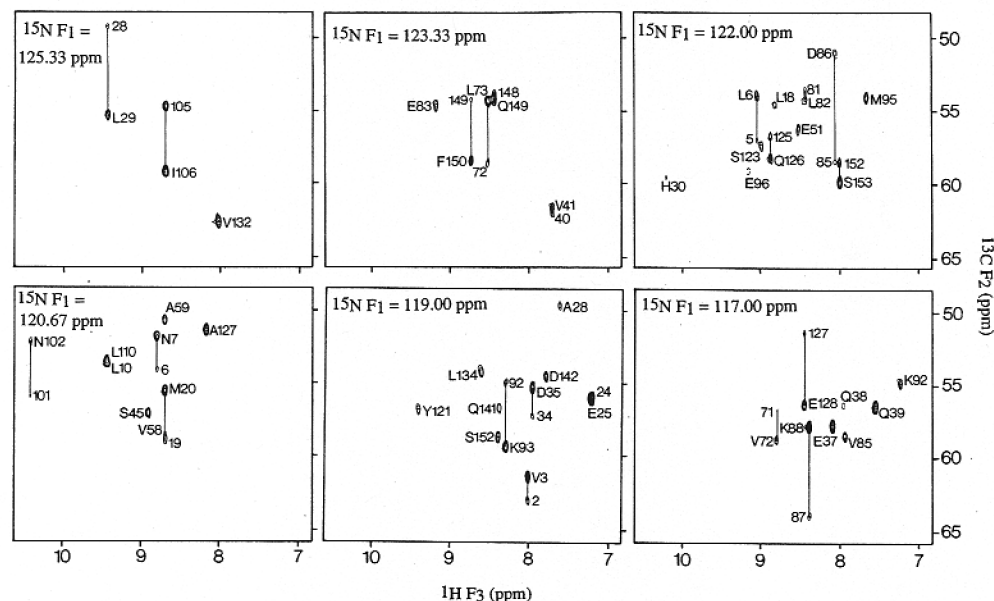


FIGURE 1: Representative slices at six different  $^{15}\text{N}(F_1)$  chemical shifts of the 3D HNCA spectrum of uniformly  $^{15}\text{N}/^{13}\text{C}$ -labeled IL-1 $\beta$ . The intrasidue  $\text{NH}(i)\text{--}^{15}\text{N}(i)\text{--}^{13}\text{C}\alpha(i)$  correlations are indicated by the residue name and number, while the weaker interresidue  $\text{NH}(i)\text{--}^{15}\text{N}(i)\text{--}^{13}\text{C}\alpha(i-1)$  cross-peaks are indicated by residue number only.

assignments for seven of the eight proline residues [see for example the Val-3(NH)–Val-3( $^{15}\text{N}$ )–Pro-2( $^{13}\text{C}\alpha$ ) cross-peak at  $\delta^{15}\text{N} = 119$  ppm in Figure 1]. From the HNCA experiment alone we were able to obtain the  $^{13}\text{C}\alpha$  chemical shifts for all but nine residues (Ala-1, Ser-17, Asn-53, Leu-62, Lys-63, Pro-91, Ser-114, Ile-143, and Thr-144) which did not show any correlations peaks. This may be due to either significant line broadening of the NH,  $^{15}\text{N}$ , or  $^{13}\text{C}\alpha$  line widths or to a reduction in NH resonance intensity arising from NOE effects. The latter is due to the selective presaturation employed to suppress the water signal and may involve either bound water or  $\text{C}^\alpha\text{H}$  protons degenerate with the water resonance or rapidly exchanging side-chain amide and guanidinium groups. In the case of Leu-62 and Lys-63, the absence of HNCA correlation peaks is in all likelihood due to the large  $^{15}\text{N}$  line widths for these two residues arising from exchange contributions to  $T_2$  (Clare et al., 1990b).

With all the  $\text{C}^\alpha\text{H}$  resonance assignments in hand, together with the large majority of  $^{13}\text{C}\alpha$  assignments obtained from the HNCA experiment, it is now possible to use the HCCH-COSY (Bax et al., 1990a) and HCCH-TOCSY (Bax et al., 1990b) experiments to delineate the amino acid spin systems. The basis of the two experiments is similar, and they make use of the well-resolved one-bond  $^1\text{H}\text{--}^{13}\text{C}$  ( $\sim 140$  Hz) and  $^{13}\text{C}\text{--}^{13}\text{C}$  (30–40 Hz)  $J$  couplings to transfer magnetization, thereby circumventing the problems associated with conventional methodologies (e.g.,  $^1\text{H}\text{--}^1\text{H}$  COSY, HOHAHA, etc.) which rely on the poorly resolved ( $<12$  Hz)  $^1\text{H}\text{--}^1\text{H}$   $J$  couplings. The scheme for the two experiments is also similar. In the first step,  $^1\text{H}$  magnetization from a proton is transferred to its directly bonded  $^{13}\text{C}$  nucleus via the  $^1J_{\text{CH}}$  coupling in an INEPT-type manner. In the second step,  $^{13}\text{C}$  magnetization is transferred to its  $^{13}\text{C}$  neighbor(s) via the  $^1J_{\text{CC}}$  coupling. In the case of the HCCH-COSY experiment, this is achieved by a  $90^\circ$   $^{13}\text{C}$  COSY mixing pulse so that magnetization is only transferred from a  $^{13}\text{C}$  nucleus to its directly bonded  $^{13}\text{C}$  neighbors; in the HCCH-TOCSY experiment, on the other hand, isotropic mixing of  $^{13}\text{C}$  spins is achieved such that both direct and multiple relayed magnetization transfer occurs along the carbon chain (Fesik et al., 1990; Bax et al., 1990b). Finally,  $^{13}\text{C}$  magnetization is transferred back to  $^1\text{H}$  via the  $^1J_{\text{CH}}$

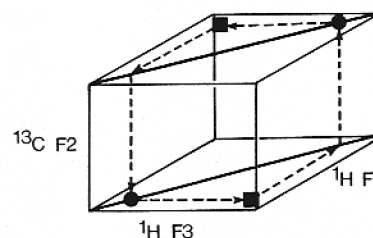


FIGURE 2: Schematic diagram of two slices of an HCCH spectrum at different  $^{13}\text{C}(F_2)$  chemical shifts showing the diagonal and cross-peaks expected for a simple two-spin system. The diagonal peaks are represented by circles and the cross-peaks by squares. Note that each slice is asymmetric about the diagonal, and cross-peaks only appear in the  $F_3$  dimension. This is due to the fact that magnetization originating from a proton attached to a given  $^{13}\text{C}$  nucleus is only visible in the  $F_2$  slice corresponding to the  $^{13}\text{C}$  chemical shift of this particular  $^{13}\text{C}$  nucleus. The corresponding cross-peak at the other side of the  $F_1 = F_3$  diagonal is found in the slice taken at the  $^{13}\text{C}$  frequency of the destination carbon.

coupling. The final result is a 3D spectrum in which each  $^1\text{H}(F_1)\text{--}^1\text{H}(F_3)$  plane has an appearance similar to that of a 2D  $^1\text{H}\text{--}^1\text{H}$  COSY or HOHAHA/TOCSY experiment but is edited by the  $^{13}\text{C}$  chemical shift of the  $^{13}\text{C}$  nucleus directly bonded to the  $^1\text{H}$  at the diagonal position from which magnetization originates. Further, in contrast to the 2D correlation experiments, the cross-peaks in each plane do not occur symmetrically on either side of the diagonal. This is represented diagrammatically in Figure 2. Consider the case where magnetization is transferred from proton A to proton B. In the plane corresponding to the  $^{13}\text{C}$  chemical shift of the  $^{13}\text{C}$  nucleus directly bonded to proton A where magnetization originates, a correlation is observed between a diagonal peak at  $(F_1, F_3) = (\delta_A, \delta_A)$  and a cross-peak at  $(F_1, F_3) = (\delta_A, \delta_B)$  in half of the spectrum. The symmetric correlation between the diagonal peak at  $(F_1, F_3) = (\delta_B, \delta_B)$  and the cross-peak at  $(F_1, F_3) = (\delta_B, \delta_A)$  is then seen in the plane corresponding to the  $^{13}\text{C}$  chemical shift of the  $^{13}\text{C}$  nucleus directly bonded to proton B. By this means, unambiguous checks on the assignments are afforded at each step in the process, which are made all the easier as the  $^{13}\text{C}$  chemical shifts for different carbon types are located in characteristic regions of the  $^{13}\text{C}$  spectrum with



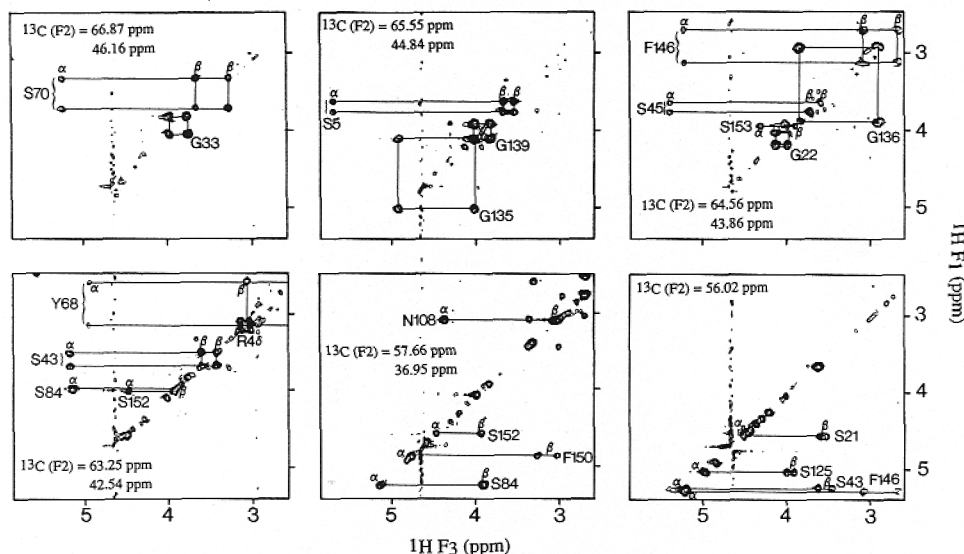


FIGURE 3: Selected slices at different  $^{13}\text{C}(F_2)$  chemical shifts of the HCCH-COSY spectrum of uniformly  $^{15}\text{N}/^{13}\text{C}$ -labeled IL-1 $\beta$ , principally illustrating connectivities involving Gly and Ser residues. Note, that owing to the small sweep width (20.71 ppm) employed in the  $^{13}\text{C}(F_2)$  dimension, extensive folding occurs, so that each slice corresponds to several  $^{13}\text{C}$  chemical shifts with a separation of 20.71 ppm. Note that the separation between successive slices in the  $F_2$  dimension is 0.32 ppm (48 Hz), so that cross-peaks originating from a particular proton may appear in adjacent slices.

little overlap between them (Howarth & Lilley, 1978). This feature also permits one to use extensive folding in the  $^{13}\text{C}(F_2)$  dimension without the risk of introducing ambiguities. Further, this enables one to resolve problems associated with extensive chemical shift degeneracy. For example, even if the  $\text{C}^\alpha\text{H}$ ,  $^{13}\text{C}\alpha$ , and  $\text{C}^\beta\text{H}$  chemical shifts of two residues are degenerate, one can still ascertain that two residues rather than one are involved and trace out the connectivities for the remainder of the two spin systems, providing the  $^{13}\text{C}\beta$  chemical shifts of the two residues are different.

The HCCH-COSY experiment is particularly useful for identifying Gly, Ala, Thr, and Val spin systems, as well as amino acids of the AMX type (e.g., Ser, Cys, Asn, Asp, His, Tyr, Phe, and Trp). This is illustrated in Figures 3–5. Figure 3 principally demonstrates the use of the HCCH-COSY experiment for the assignment of Gly and Ser spin systems. Both these residues tend to pose problems in conventional  $^1\text{H}$ – $^1\text{H}$  correlation spectra as the proton resonances are often very close to each other. Each slice corresponds to several  $^{13}\text{C}$  shifts separated by 20.71 ppm. Gly  $^{13}\text{C}\alpha$  shifts occur at 42–46 ppm, while  $^{13}\text{C}\alpha$  and  $^{13}\text{C}\beta$  shifts of Ser occur at 55–60 and 61–67 ppm, respectively. The two Gly  $\text{C}^\alpha\text{H}$  protons are attached to the same  $\text{C}\alpha$  carbon, so that a symmetric pattern appears about the diagonal. The same is true, of course, for  $\beta$ -methylene protons. However, the  $\beta$ -methylene protons are also associated with cross-peaks to the  $\text{C}^\alpha\text{H}$  proton. Consequently, there is no difficulty in deciding which peaks originate from  $\text{C}^\alpha\text{H}$  protons of Gly and  $\text{C}^\beta\text{H}$  protons of either Ser (left-hand and middle top panels of Figure 3) or another AMX spin system such as Phe whose  $\text{C}\beta$  has a chemical shift similar to that of the Gly  $\text{C}\alpha$  (right-hand top panel of Figure 3). The assignment of the Ser spin system is easily checked by examining the slices corresponding to the  $\text{C}\alpha$  and  $\text{C}\beta$  chemical shifts of Ser (e.g., left-hand and middle bottom panels of Figure 3). In this respect, we note that in conventional  $^1\text{H}$ – $^1\text{H}$  correlated spectra it is often difficult to ascertain the  $\text{C}^\alpha\text{H}$  and  $\text{C}^\beta\text{H}$  chemical shifts of Ser for cases where the  $\text{C}^\beta\text{H}$  shifts are degenerate, as the latter may lie to both low and high field of the  $\text{C}^\alpha\text{H}$  shifts. In the case of the HCCH-COSY experiment, this is simply resolved as the  $^{13}\text{C}$  chemical shifts of the  $\text{C}\alpha$  and  $\text{C}\beta$  carbons differ by 2–10 ppm, with the  $\text{C}\alpha$  resonance

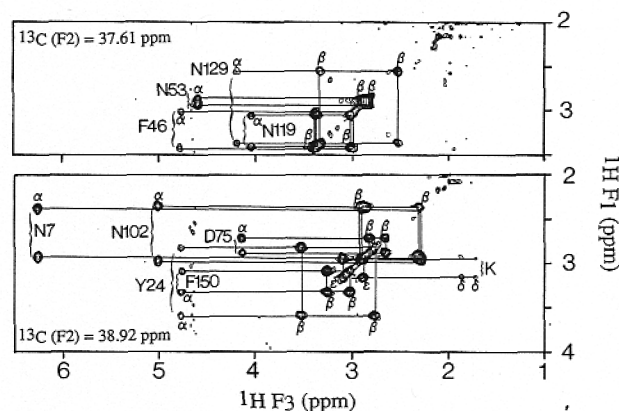


FIGURE 4: Selected slices at different  $^{13}\text{C}(F_2)$  chemical shifts of the HCCH-COSY spectrum of uniformly  $^{15}\text{N}/^{13}\text{C}$ -labeled IL-1 $\beta$  illustrating connectivities from the  $\text{C}^\beta\text{H}$  to  $\text{C}^\alpha\text{H}$  protons for AMX spin systems. In the bottom panel, two cross-peaks from the  $\text{C}^\beta\text{H}$  to the  $\text{C}^\alpha\text{H}$  protons of a lysine residue are assigned. These peaks originate from one of the four lysine residues (out of a total of 15 lysines present in IL-1 $\beta$ ) for which complete side-chain assignments could not be obtained.

always upfield of the  $\text{C}\beta$  one. Examples of other cross-peaks arising from the  $\text{C}^\beta\text{H}$  protons of AMX spin systems are illustrated in Figure 4 for a number of Tyr, Phe, Asn, and Asp residues. The assignment of Ala spin systems is also completely straightforward as only two  $^{13}\text{C}$  and  $^1\text{H}$  shifts are involved.

Thr and Val spin systems display characteristic connectivities originating from the  $\text{C}^\beta\text{H}$  protons to the  $\text{C}^\alpha\text{H}$  and  $\text{C}^\gamma\text{H}$  methyl protons. The Thr  $\text{C}\beta$  chemical shifts are at 70–73 ppm, while those of Val are at 31–37 ppm, so the cross-peaks arising from a  $\text{C}^\beta\text{H}$  proton of a Val or Thr residue can be distinguished (irrespective of arguments based on  $\text{C}^\beta\text{H}$  chemical shifts) even when the chemical shifts of the two methyl proton resonances of Val are degenerate. Some examples of connectivities arising from the  $\text{C}^\beta\text{H}$  proton of Thr are illustrated in Figure 6 (middle panel).

Despite the resolution afforded by the  $^{13}\text{C}(F_2)$  dimension in the HCCH-COSY experiment, problems still remain with regard to the assignment of other spin systems. For example,

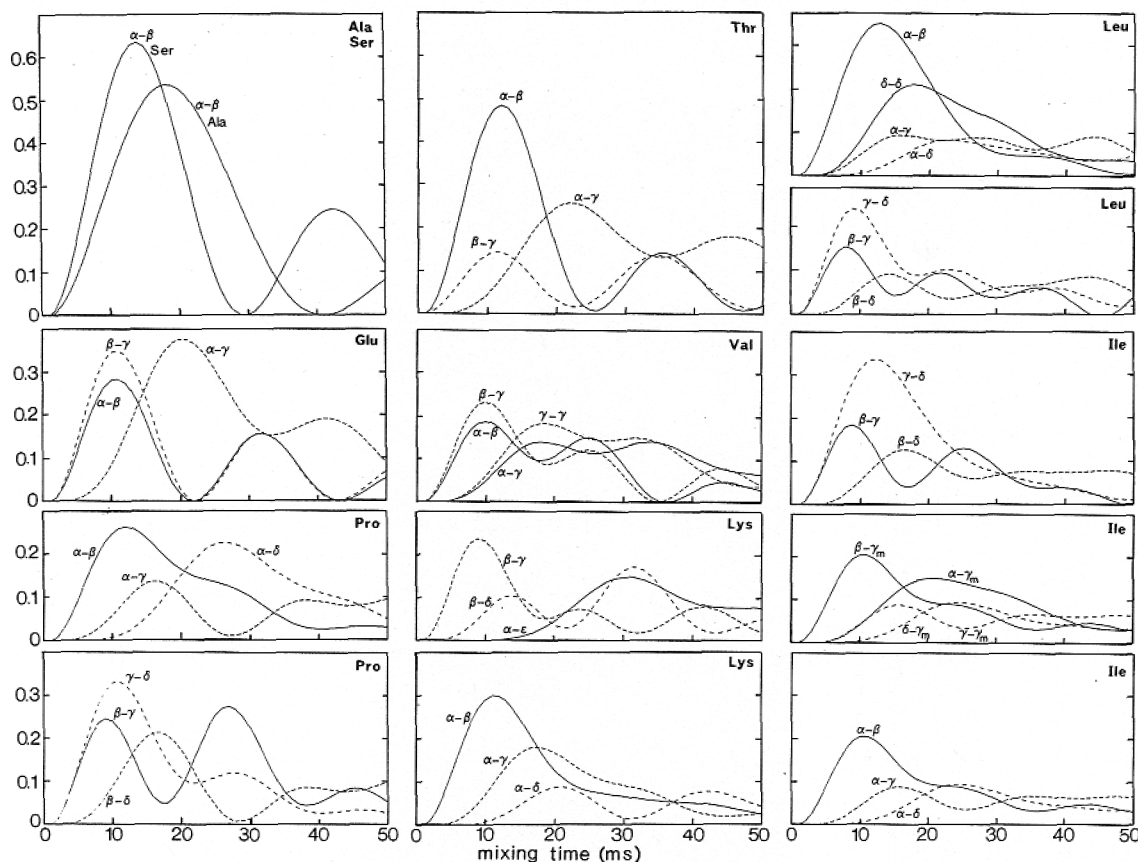


FIGURE 5: Amount of net intraresidue carbon-carbon magnetization transfer occurring during the DIPSI-3 mixing period in the HCCH-TOCSY experiment, assuming a relaxation rate of  $^{13}\text{C}$  magnetization during isotropic mixing equal to  $33\text{ s}^{-1}$ . Because of  $J_{\text{CC}}$  dephasing during the  $\delta 1$ ,  $\delta 2$ , and  $t_2$  periods in the pulse sequence (see Experimental Procedures), the curves should be shifted toward the left by several milliseconds for predicting the cross-peak intensities observed in the HCCH-TOCSY spectrum. Cross-peak intensities in the HCCH-TOCSY experiment are proportional to the amount of carbon-carbon magnetization transfer but are attenuated 2-fold if one of the carbons has two magnetically inequivalent protons attached and 4-fold if both carbons have nonequivalent protons attached. Reduced effective carbon-carbons  $J$  couplings (Bax et al., 1990b) have been used for calculating these graphs: Ser,  $J_{\alpha\beta} = 35\text{ Hz}$ ; Ala,  $J_{\alpha\beta} = 25\text{ Hz}$ ; Thr,  $J_{\alpha\beta} = 20\text{ Hz}$ ,  $J_{\beta\gamma} = 35\text{ Hz}$ ; Glu,  $J_{\alpha\beta} = 30\text{ Hz}$ ,  $J_{\beta\gamma} = 33\text{ Hz}$ ; Val,  $J_{\alpha\beta} = 27\text{ Hz}$ ,  $J_{\beta\gamma} = 30\text{ Hz}$ ; Pro,  $J_{\alpha\beta} = 27\text{ Hz}$ ,  $J_{\beta\gamma} = 33\text{ Hz}$ ; Lys,  $J_{\alpha\beta} = 30\text{ Hz}$ ,  $J_{\beta\gamma} = J_{\beta\delta} = J_{\delta\epsilon} = 33\text{ Hz}$ ; Leu,  $J_{\alpha\beta} = J_{\beta\gamma} = 30\text{ Hz}$ ,  $J_{\gamma\delta} = 33\text{ Hz}$ ; Ile,  $J_{\alpha\beta} = J_{\beta\gamma} = 27\text{ Hz}$ ,  $J_{\beta\gamma} = J_{\gamma\delta} = 30\text{ Hz}$ . All multiple-bond carbon-carbon couplings are assumed to be zero, and couplings between aliphatic and carbonyl resonances have been neglected because the large resonance chemical shift difference between the aromatic and carbonyl nuclei on the one hand and the aliphatic carbons on the other reduces their effective  $J$  couplings to aliphatic carbons (Shaka et al., 1988; Bax et al., 1990b).

it is frequently the case for Gln, Glu, Met, and Leu spin systems that the chemical shifts of the  $\text{C}^\beta\text{H}$  and  $\text{C}^\gamma\text{H}$  protons are similar and sometimes degenerate. Consequently, it is often difficult to locate the  $\text{C}^\beta\text{H}$ - $\text{C}^\gamma\text{H}$  connectivities unambiguously in the HCCH-COSY spectrum. These problems are even further exacerbated for longer side-chain spin systems such as Arg, Pro, and Lys. The same problem presents itself in the interpretation of regular 2D  $^1\text{H}$ - $^1\text{H}$  COSY type spectra and can be circumvented by means of experiments that detect not only direct but relayed through-bond connectivities, for example, relayed COSY and HOHAHA/TOCSY spectroscopy. Similarly, the ambiguities present for longer side-chain spin systems in the HCCH-COSY experiment can be resolved unambiguously by recording an HCCH-TOCSY experiment in which carbon magnetization is transferred along the carbon chain by isotropic mixing.

In order to ensure that the HCCH-TOCSY experiment is recorded under conditions which are optimal for the observation of relayed connectivities, it is essential to consider in detail the expected theoretical cross-peak intensities as a function of isotropic mixing time. As discussed previously (Bax et al., 1990b), the DIPSI-3 sequence (Shaka et al., 1988) is ideally suited for isotropic mixing of  $^{13}\text{C}$  magnetization. For the case of  $^{13}\text{C}$  with relatively large homonuclear couplings (30–40 Hz), the DIPSI-3 sequence can cover a bandwidth of

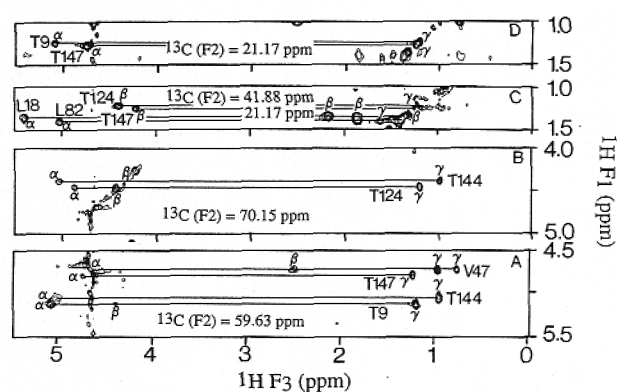


FIGURE 6: Selected slices at different  $^{13}\text{C}(F_2)$  chemical shifts of the HCCH-TOCSY (first and fourth panels) and HCCH-COSY (second and third panels) spectra of uniformly  $^{15}\text{N}/^{13}\text{C}$ -labeled IL-1 $\beta$  principally illustrating Thr spin systems. Note that the direct  $\text{C}^\gamma\text{H}$ - $\text{C}^\beta\text{H}$  and  $\text{C}^\alpha\text{H}$ - $\text{C}^\beta\text{H}$  cross-peaks for the threonine spin systems are either very weak or absent in the HCCH-TOCSY spectrum but are strong in the HCCH-COSY spectrum.

ca.  $\pm 0.75\nu$ , where  $\nu$  is the strength of the applied  $^{13}\text{C}$  rf field. Thus, on a 500-MHz spectrometer, an rf field strength of only  $\sim 7\text{ kHz}$  is needed for efficient isotropic mixing of the aliphatic resonances. On our inverse  $^1\text{H}$  probe head, this requires 12.5 W of rf power, which is sufficiently low to permit mixing for

durations as long as 25 ms without causing substantial sample heating. The rate at which magnetization is transferred from one carbon to its neighbor depends on the size of the effective  $J$  coupling,  $J_{\text{eff}}$ , present during the mixing period. For coupled  $^{13}\text{C}$  nuclei with widely different chemical shifts (e.g., Thr- $\text{C}\beta$  and  $-\text{C}\gamma$ ), the effective  $J$  coupling can be substantially smaller than the real  $J$  coupling (Bax et al., 1990b). This results in a significant reduction in the rate of  $^{13}\text{C}$  magnetization transfer which must be taken into account in simulations of the magnetization transfer process.

Figure 5 shows the calculated fraction of magnetization that is transferred from one carbon to another as a function of the length of the isotropic mixing period for various amino acids. These simulations employ approximate values for the effective  $J$  couplings that are smaller than the real  $J$  couplings (Tran-Dinh et al., 1974; Jogn et al., 1974; Krivdin & Kalabin, 1989). This apparent reduction depends on the difference between the  $^{13}\text{C}$  shifts of the coupled  $^{13}\text{C}$  resonances and was estimated with an rf field of 7 kHz and a  $^{13}\text{C}$  frequency of 125.76 MHz (i.e., a 500-MHz spectrometer). It should be noted, however, that the general shape of these curves does not depend strongly on either the rf field or the magnetic field strength. In all the simulations, the presence of carbonyl or aromatic carbons has been neglected because the effective  $J$  coupling during isotropic mixing to the aliphatic carbon is very small (<10 Hz). In addition, for simplicity, the magnetization of all  $^{13}\text{C}$  nuclei was assumed to relax at a rate of  $33\text{ s}^{-1}$  during the isotropic mixing period. Because in-phase and antiphase magnetization components that participate in the magnetization transfer process relax at different rates, a proper calculation of the pertinent relaxation rates occurring during isotropic mixing is rather complex. The relaxation rate used for the simulation corresponds to the magnetization transfer pathway that results in cross-peaks, assuming a chain of methine carbons, tumbling isotropically with a correlation time of 10 ns and an order parameter  $S^2$  of 0.85. In practice, of course,  $^{13}\text{C}$  relaxation rates will vary over a substantial range, depending on the local order parameter (i.e., internal mobility) and on the number of protons attached to the  $^{13}\text{C}$  nucleus. Thus, the curves shown in Figure 5 provide only an approximate, but very useful, guide of what magnetization transfer to expect for a given mixing time duration. Calculation of exact curves would require accurate knowledge of  $^1J_{\text{CC}}$  couplings,  $^{13}\text{C}$  chemical shifts, and longitudinal and transverse  $^{13}\text{C}$  relaxation rates. Cross-peak intensities observed in the HCCH-TOCSY experiment are proportional to the amount of carbon-carbon magnetization transfer but are attenuated 2-fold if one of the carbons has two magnetically inequivalent protons attached and 4-fold if both carbons have nonequivalent protons attached.

For the simplest case of a two-spin system, transfer of magnetization from one carbon to its neighbor occurs in a cosinusoidal manner (Braunschweiler & Ernst, 1983), with a period of  $1/J_{\text{eff}}$ . For Ala, the  $\text{C}\alpha$  and  $\text{C}\beta$  resonances differ in chemical shift by  $\sim 33$  ppm, and consequently,  $J_{\text{eff}}$  (25 Hz) is much smaller than the real  $J_{\text{C}\alpha\text{C}\beta}$  coupling (35 Hz). For Ser, the difference in  $\text{C}\alpha$  and  $\text{C}\beta$  shifts is typically small (<10 ppm) so that the effective  $J$  coupling will be close to its normal value (37 Hz). As a result, magnetization transfer between the  $\text{C}\alpha$  and  $\text{C}\beta$  carbons occurs faster for Ser than for Ala residues (top left-hand panel in Figure 5). For other two-carbon systems (i.e., His, Phe, Tyr, Trp, Cys, Asp, and Asn) magnetization occurs at a rate that is intermediate between the Ala and Ser curves. Similarly, the  $J_{\text{C}\alpha\text{C}\beta}$  coupling for Thr is close to its regular value as the  $\text{C}\alpha - \text{C}\beta$  chemical shift

difference is typically less than 10 ppm, whereas the  $J_{\text{C}\beta\text{C}\gamma}$  coupling is reduced by almost 50% because of the large  $\text{C}\beta - \text{C}\gamma$  chemical shift difference ( $\sim 50$  ppm). Even though the carbon topology of the Thr side chain is the same as that for Glu, the reduction in the Thr  $J_{\text{C}\beta\text{C}\gamma}$  coupling caused by offset effects makes the transfer of magnetization from the  $\text{C}\beta$  to the  $\text{C}\gamma$  carbon considerably less efficient in Thr than in Glu, as is readily seen in Figure 5. The transfer efficiency shown for Glu also applies to Gln and Met.

The magnetization transfer efficiencies in Pro and Arg are also very similar, although transfer to the  $\text{C}\delta$  of Arg may often appear more intense than that for the  $\text{C}\delta$  of Pro due to the frequently observed increase in mobility at the end of the long Arg side chain. For both residues the effective  $J_{\text{C}\alpha\text{C}\beta}$  coupling is significantly smaller than the real  $J$  coupling because of the large difference in  $\text{C}\alpha$  and  $\text{C}\beta$  chemical shifts. Note that if all one-bond  $J_{\text{CC}}$  couplings were identical, transfer from  $\text{C}\alpha$  to  $\text{C}\beta$  would follow the same curve as transfer from  $\text{C}\delta$  to  $\text{C}\gamma$  and transfer from  $\text{C}\alpha$  to  $\text{C}\gamma$  would be identical with transfer from  $\text{C}\delta$  to  $\text{C}\beta$ . For Lys,  $J_{\text{C}\alpha\text{C}\beta}$  is expected to be only slightly smaller than the other one-bond couplings. Thus, transfer from  $\text{C}\epsilon$  to  $\text{C}\delta$  (not shown) follows the same pattern as from  $\text{C}\alpha$  to  $\text{C}\beta$ , and so on.

The curves shown in Figure 5 are calculated for net magnetization transfer during the isotropic mixing period alone. In practice, however, dephasing of  $^{13}\text{C}$  magnetization also occurs during the  $\delta 1$  and  $\delta 2$  periods (cf. the HCCH-TOCSY pulse scheme under Experimental Procedures). This has the effect of shifting the curves displayed in Figure 5 to the left by several milliseconds. In addition,  $J_{\text{CC}}$  dephasing occurs during the evolution period  $t_2$ . This also results in a left shift of the curves, but by a variable amount depending on the  $t_2$  duration, and, in addition, causes  $F_2$  line shapes to deviate slightly from the ideal absorptive phase.

From analysis of the theoretical curves shown in Figure 5, we concluded that an isotropic mixing time in the range 20–30 ms was optimal for the observation of relayed connectivities. Examples of different sorts of relayed connectivities observed in the 24-ms mixing time HCCH-TOCSY spectra are illustrated in Figures 6–10.

Figure 6 illustrates some Thr spin systems and affords a comparison of the HCCH-COSY and HCCH-TOCSY spectra. The top two panels in Figure 6 are taken at the same  $^{13}\text{C}$  chemical shift of 21.17 ppm and illustrate the connectivities that are observed from the  $\text{C}^{\gamma}\text{H}$  proton of Thr in the 24-ms HCCH-TOCSY (first panel) and the HCCH-COSY (second panel) spectra. In the HCCH-COSY spectrum, a strong cross-peak corresponding to the direct connectivity to the  $\text{C}^{\beta}\text{H}$  proton is observed, while in the HCCH-TOCSY spectrum a strong cross-peak arises from the relayed  $\text{C}^{\gamma}\text{H}-\text{C}^{\alpha}\text{H}$  connectivity. As predicted from the simulations (Figure 5), the direct  $\text{C}^{\gamma}\text{H}-\text{C}^{\beta}\text{H}$  connectivity is not observed in the HCCH-TOCSY spectrum recorded with a 24-ms isotropic mixing time. In the case of Val, on the other hand, the intensities of the relayed  $\text{C}^{\gamma}\text{H}-\text{C}^{\alpha}\text{H}$  and direct  $\text{C}^{\gamma}\text{H}-\text{C}^{\beta}\text{H}$  cross-peaks are comparable and slightly weaker than that of the relayed  $\text{C}^{\gamma}\text{H}-\text{C}^{\beta}\text{H}$  cross-peak between the two methyl carbons (cf. the theoretical curves in Figure 5 and the experimental data in the bottom panel of Figure 6 and the top two panels of Figure 9).

Figure 7 illustrates relayed connectivities originating from the  $\text{C}^{\alpha}\text{H}$  proton for a number of different amino acid spin systems in the 24-ms HCCH-TOCSY spectrum. Note that complete spin systems may be delineated not only for AM-(PT)X spin systems such as Glu, Gln, and Met (see slices at 53.38, 55.03, 56.68, 57.99, and 58.65 ppm) but also for Val,



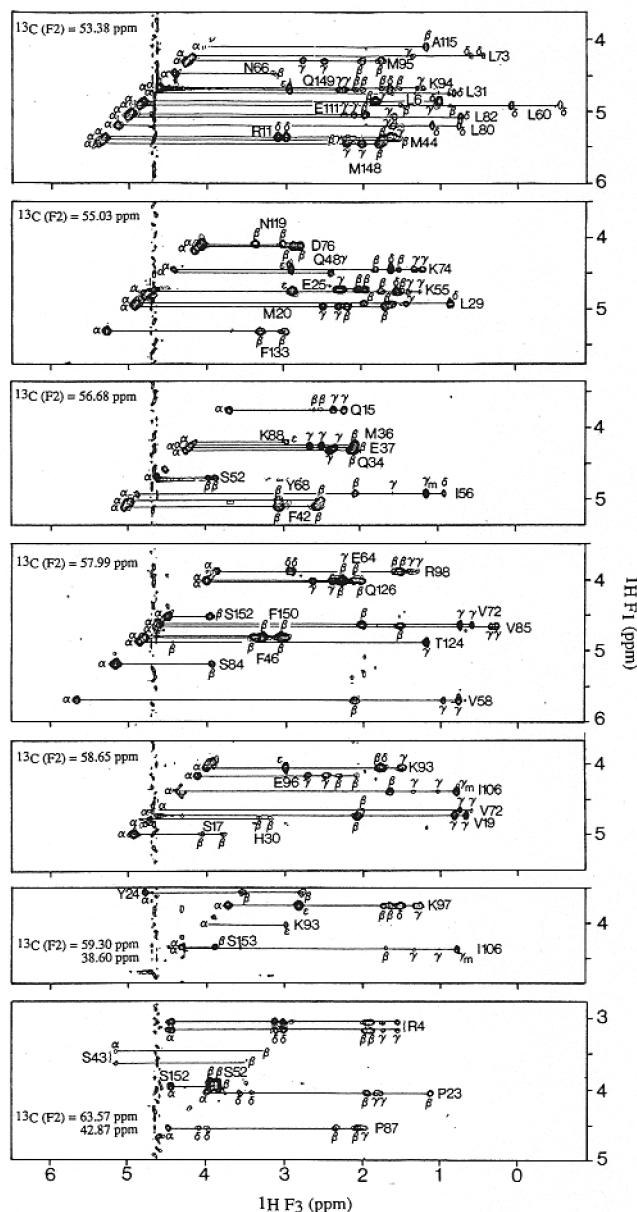


FIGURE 7: Selected slices at different  $^{13}\text{C}(F_2)$  chemical shifts of the HCCH-TOCSY spectrum of uniformly  $^{15}\text{N}/^{13}\text{C}$ -labeled IL-1 $\beta$  illustrating relayed connectivities originating from the  $\text{C}^\alpha\text{H}$  proton of several longer side chain amino acid spin systems such as Glu, Gln, Met, Pro, Arg, and Lys.

Pro, Arg, and Lys. It should be noted, however, that the  $\text{C}^\beta$  and  $\text{C}^\gamma$  chemical shifts of Met are very similar, and likewise, the difference in  $\text{C}^\beta$  and  $\text{C}^\gamma$  chemical shift for Gln and Glu may be as small as 1 ppm (e.g., Glu-111 in Table I). Consequently, it is essential in these cases to verify the identification of the direct  $\text{C}^\alpha\text{H}-\text{C}^\beta\text{H}$  connectivities with the HCCH-COSY experiment, even though the intensity of the  $\text{C}^\alpha\text{H}-\text{C}^\gamma\text{H}$  cross-peaks is usually significantly stronger, and never weaker, than those of the direct  $\text{C}^\alpha\text{H}-\text{C}^\beta\text{H}$  cross-peaks in the 24-ms HCCH-TOCSY spectrum. Examples of Val spin systems are evident at  $\delta^{13}\text{C} = 57.99$  ppm (Val-58 and Val-85) and  $\delta^{13}\text{C} = 58.65$  ppm (Val-19 and Val-72). In the case of the Val spin systems, the intensities of the direct  $\text{C}^\alpha\text{H}-\text{C}^\beta\text{H}$  and relayed  $\text{C}^\alpha\text{H}-\text{C}^\gamma\text{H}$  cross-peaks are comparable, although the latter tend to be slightly stronger because the  $\text{C}^\gamma\text{H}$  proton resonance is narrower than the  $\text{C}^\beta\text{H}$  one. Complete spin systems for several Leu residues (Leu-6, Leu-29, Leu-31, Leu-60, Leu-73, Leu-80, and Leu-82) may be seen in the slices at  $\delta^{13}\text{C} = 53.38$  and 55.03 ppm, and the intensities of the direct  $\text{C}^\alpha\text{H}-\text{C}^\beta\text{H}$  and

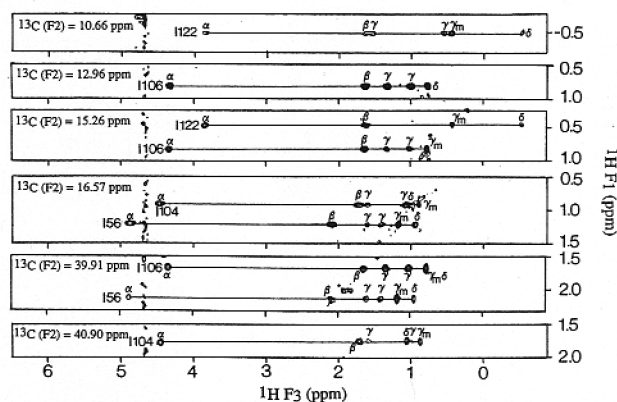


FIGURE 8: Selected slices at different  $^{13}\text{C}(F_2)$  chemical shifts of the HCCH-TOCSY spectrum of uniformly  $^{15}\text{N}/^{13}\text{C}$ -labeled IL-1 $\beta$  illustrating multiple relayed connectivities originating from the  $\text{C}^\alpha\text{H}_3$ ,  $\text{C}^\gamma\text{H}_3$ , and  $\text{C}^\beta\text{H}$  protons of several Ile spin systems.

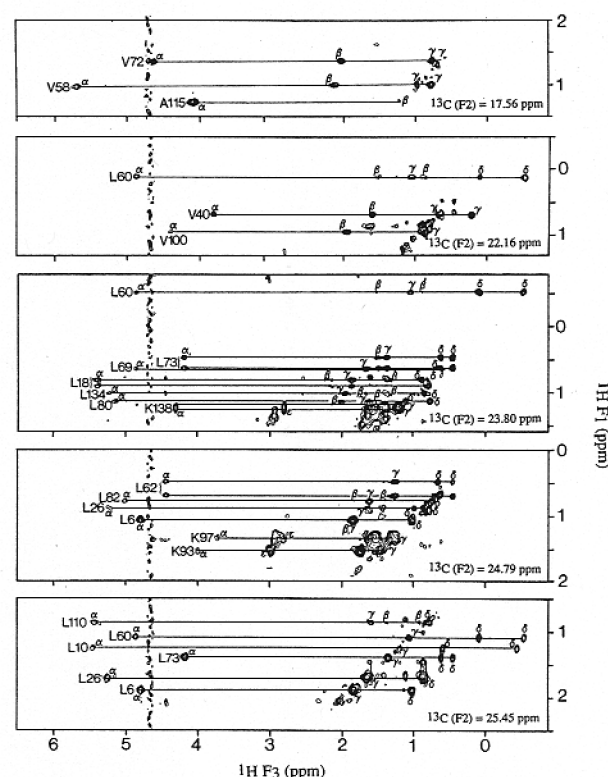


FIGURE 9: Selected slices at different  $^{13}\text{C}(F_2)$  chemical shifts of the HCCH-TOCSY spectrum of uniformly  $^{15}\text{N}/^{13}\text{C}$ -labeled IL-1 $\beta$  illustrating multiple relayed connectivities originating from the  $\text{C}^\gamma\text{H}_3$  protons of Val, the  $\text{C}^\alpha\text{H}_3$  and  $\text{C}^\gamma\text{H}$  protons of Leu, and the  $\text{C}^\gamma\text{H}$  protons of Lys.

relayed  $\text{C}^\alpha\text{H}-\text{C}^\gamma\text{H}$  and  $\text{C}^\alpha\text{H}-\text{C}^\beta\text{H}$  cross-peaks are usually comparable, as expected for a mixing time of 24 ms from the theoretical curves in Figure 5. Two Ile residues at  $\delta^{13}\text{C} = 56.68$  ppm (Ile-56) and  $\delta^{13}\text{C} = 58.65$ –59.30 ppm (Ile-106; note that the  $\text{C}^\gamma\text{H}$  and  $\text{C}^\beta\text{H}$  methyl resonances are superimposed) are also present in Figure 7. For this spin system, the ratio of cross-peak intensities tends to follow the pattern  $\text{C}^\alpha\text{H}-\text{C}^\gamma\text{H} > \text{C}^\alpha\text{H}-\text{C}^\beta\text{H} \sim \text{C}^\alpha\text{H}-\text{C}^\beta\text{H} > \text{C}^\alpha\text{H}-\text{C}^\gamma\text{H}$ . Although the latter cross-peaks are weak, they are still clearly detectable, which is rarely the case in conventional  $^1\text{H}-^1\text{H}$  2D HOHAHA spectra of much smaller proteins of 40–60 residues.

As discussed above in relation to the theoretical simulations, the behavior of Pro and Arg spin systems in the HCCH-TOCSY spectrum should be similar, with the proviso that connectivities to the  $\text{C}^\beta\text{H}$  proton of arginine would be expected

Gly									Q38				Q39				Q48			
<sup>1</sup> H	G22 Ca		G33 Ca		G49 Ca		G61 Ca		<sup>1</sup> H	Ca	Cβ	Cγ	Ca	Cβ	Cγ	Ca	Cβ	Cγ		
<sup>13</sup> C	4.12,3.97		3.99,3.77		4.12,3.94		4.15,2.38		<sup>13</sup> C	4.21	2.23,1.97	2.50,2.40	4.42	2.30,2.03	2.56,2.42	4.44	2.08,1.97	2.42,2.42		
	43.85		46.33		44.51		42.05			55.52	29.06	33.67	55.85	29.22	33.99	55.36	29.06	33.34		
										E50			E51			E64				
<sup>1</sup> H	G135 Ca		G136 Ca		G139 Ca		G140 Ca		<sup>1</sup> H	Ca	Cβ	Cγ	Ca	Cβ	Cγ	Ca	Cβ	Cγ		
<sup>13</sup> C	4.92,4.02		3.84,2.93		4.02,3.83		4.20,3.78		<sup>13</sup> C	4.35	2.06,1.92	2.25,2.25	4.50	2.11,2.01	2.34,2.27	3.98	2.24,2.09	2.24,2.24		
Ala	44.68		44.02		45.00		44.35			55.85	29.22	35.64	55.52	30.21	35.64	57.82	27.91	33.34		
										E83			E83			E96				
										E105			E111			E113				
<sup>1</sup> H	A1		A28		A115		A59		<sup>1</sup> H	Ca	Cβ	Cγ	Ca	Cβ	Cγ	Ca	Cβ	Cγ		
<sup>13</sup> C	4.36		6.16		4.05		5.42		<sup>13</sup> C	4.98	2.11,2.11	2.69,2.51	4.63	2.10,1.95	2.31,2.31	4.09	2.23,2.08	2.73,2.44		
	50.26		49.95		53.06		50.26			53.06	31.69	32.18	54.05	32.84	36.62	58.65	28.08	36.95		
										E105			E111			E113				
<sup>1</sup> H	A127								<sup>1</sup> H	Ca	Cβ	Cγ	Ca	Cβ	Cγ	Ca	Cβ	Cγ		
<sup>13</sup> C	4.38								<sup>13</sup> C	5.14	1.94,1.73	2.14,2.05	5.00	1.99,1.99	2.22,2.11	4.47	1.42,1.42	2.66,2.56		
Ser	50.92		19.20							54.21	30.54	35.47	53.06	31.52	32.35	53.87	28.73	33.99		
										Q116			Q126			E128				
										Ca	Cβ	Cγ	Ca	Cβ	Cγ	Ca	Cβ	Cγ		
<sup>1</sup> H	S5		S13		S17		S21		<sup>1</sup> H	3.61	1.42,0.98	1.31,0.72	3.98	2.24,2.01	2.32,2.38	4.04	2.27,1.96	2.56,2.50		
<sup>13</sup> C	5.67		3.99		4.91		4.52		<sup>13</sup> C	56.02	28.73	32.35	57.66	28.24	33.34	53.52	30.71	34.52		
	56.34		59.79		58.47		56.18			Q141			Q149							
	S43		S45		S52		S70		<sup>1</sup> H	Ca	Cβ	Cγ	Ca	Cβ	Cγ	Ca	Cβ	Cγ		
<sup>1</sup> H	5.16		3.99		4.63		5.23		<sup>13</sup> C	4.29	2.20,2.00	2.39,2.39	4.63	2.07,1.99	2.31,2.25					
<sup>13</sup> C	55.85		56.35		56.84		55.19		Thr	55.85	28.89	33.01	53.55	29.55	33.67					
	S84		S114		S123		S125		<sup>1</sup> H	Ca	Cβ	Cγ	Ca	Cβ	Cγ	Ca	Cβ	Cγ		
<sup>1</sup> H	5.15		3.90		5.82		4.97		<sup>13</sup> C	5.06	4.39	1.18	4.58	3.93	1.20	4.83	4.39	1.16		
<sup>13</sup> C	57.82		58.65		57.00		56.18			59.30	71.47	21.01	61.60	72.28	20.35	57.99	70.31	21.01		
	S152		S153						<sup>1</sup> H	Ca	Cβ	Cγ	Ca	Cβ	Cγ	Ca	Cβ	Cγ		
<sup>1</sup> H	4.48		4.31		5.01		4.27		<sup>13</sup> C	4.31	3.99	0.96	4.99	4.39	0.96	4.74	4.19	1.22		
Asn/Asp	57.82		59.14																	

Table I (Continued)

Arg					R11					K77				
$^1\text{H}$	$\text{C}\alpha$	$\text{C}\beta$	$\text{C}\gamma$	$\text{C}\delta$	$\text{C}\alpha$	$\text{C}\beta$	$\text{C}\gamma$	$\text{C}\delta$	$\text{C}\epsilon$	$^1\text{H}$	$\text{C}\alpha$	$\text{C}\beta$	$\text{C}\gamma$	$\text{C}\delta$
$^{13}\text{C}$	4.50	2.00,1.94	1.77,1.58	3.17,3.06	5.31	1.64,1.64	1.67,1.67	3.08,3.00		$^{13}\text{C}$	4.66	1.67,1.60	1.39,1.39	2.86
	55.69	30.71	27.09	42.70	53.54	34.15	26.92	43.36			52.56	33.99	23.31	41.23
R98					R11					K88				
$^1\text{H}$	$\text{C}\alpha$	$\text{C}\beta$	$\text{C}\gamma$	$\text{C}\delta$	$\text{C}\alpha$	$\text{C}\beta$	$\text{C}\gamma$	$\text{C}\delta$	$\text{C}\epsilon$	$^1\text{H}$	$\text{C}\alpha$	$\text{C}\beta$	$\text{C}\gamma$	$\text{C}\delta$
$^{13}\text{C}$	3.89	1.42,1.33	1.58,1.53	2.97,2.92	5.31	1.64,1.64	1.67,1.67	3.08,3.00		$^{13}\text{C}$	4.24	1.83,1.72	1.47,1.47	2.84
	58.15	32.35	29.06	41.23	53.54	34.15	26.92	43.36			57.16	31.52	24.29	41.06
Lys					K92					K93				
$^1\text{H}$	$\text{C}\alpha$	$\text{C}\beta$	$\text{C}\gamma$	$\text{C}\delta$	$\text{C}\alpha$	$\text{C}\beta$	$\text{C}\gamma$	$\text{C}\delta$	$\text{C}\epsilon$	$^1\text{H}$	$\text{C}\alpha$	$\text{C}\beta$	$\text{C}\gamma$	$\text{C}\delta$
$^{13}\text{C}$	3.66	2.00,0.93	1.24,1.38	2.00,1.24	2.65,2.65	41.23				$^{13}\text{C}$	4.00	1.75,1.75	1.53,1.47	2.86
	58.32	29.92									58.80	32.18	24.62	41.06
K16					K94					K97				
$^1\text{H}$	$\text{C}\alpha$	$\text{C}\beta$	$\text{C}\gamma$	$\text{C}\delta$	$\text{C}\alpha$	$\text{C}\beta$	$\text{C}\gamma$	$\text{C}\delta$	$\text{C}\epsilon$	$^1\text{H}$	$\text{C}\alpha$	$\text{C}\beta$	$\text{C}\gamma$	$\text{C}\delta$
$^{13}\text{C}$	5.17	1.89,1.49	1.19,1.00	1.49,1.49	2.78,2.78	41.71				$^{13}\text{C}$	4.61	1.78,1.56	1.33,1.25	2.86
	54.53	36.95	24.86	29.06	41.71						53.72	33.99	24.46	41.71
K55					K99					K103				
$^1\text{H}$	$\text{C}\alpha$	$\text{C}\beta$	$\text{C}\gamma$	$\text{C}\delta$	$\text{C}\alpha$	$\text{C}\beta$	$\text{C}\gamma$	$\text{C}\delta$	$\text{C}\epsilon$	$^1\text{H}$	$\text{C}\alpha$	$\text{C}\beta$	$\text{C}\gamma$	$\text{C}\delta$
$^{13}\text{C}$	4.72	1.78,1.50	1.44,1.27	1.58,1.58	2.92,2.92	41.23				$^{13}\text{C}$	3.75	1.75,1.65	1.53,1.53	2.86
	54.86	33.99	24.29	29.06	41.23						59.13	31.85	24.61	41.06
K63					K109					K138				
$^1\text{H}$	$\text{C}\alpha$	$\text{C}\beta$	$\text{C}\gamma$	$\text{C}\delta$	$\text{C}\alpha$	$\text{C}\beta$	$\text{C}\gamma$	$\text{C}\delta$	$\text{C}\epsilon$	$^1\text{H}$	$\text{C}\alpha$	$\text{C}\beta$	$\text{C}\gamma$	$\text{C}\delta$
$^{13}\text{C}$	3.84	-	-	-	-	-	-	-	-	$^{13}\text{C}$	4.64	-	-	-
	-	-	-	-	-	-	-	-	-		55.52	-	-	-
K65					K138					K138				
$^1\text{H}$	$\text{C}\alpha$	$\text{C}\beta$	$\text{C}\gamma$	$\text{C}\delta$	$\text{C}\alpha$	$\text{C}\beta$	$\text{C}\gamma$	$\text{C}\delta$	$\text{C}\epsilon$	$^1\text{H}$	$\text{C}\alpha$	$\text{C}\beta$	$\text{C}\gamma$	$\text{C}\delta$
$^{13}\text{C}$	4.48	1.77,1.42	1.00,0.81	1.42,1.36	3.00,3.00	41.39				$^{13}\text{C}$	4.81	1.94,1.86	1.47,1.39	2.86
	54.21	34.32	23.80	29.06	41.39						54.15	34.71	24.46	41.23
K74					K138					K138				
$^1\text{H}$	$\text{C}\alpha$	$\text{C}\beta$	$\text{C}\gamma$	$\text{C}\delta$	$\text{C}\alpha$	$\text{C}\beta$	$\text{C}\gamma$	$\text{C}\delta$	$\text{C}\epsilon$	$^1\text{H}$	$\text{C}\alpha$	$\text{C}\beta$	$\text{C}\gamma$	$\text{C}\delta$
$^{13}\text{C}$	4.42	1.86,1.56	1.33,1.22	1.67,1.67	2.94,2.94	40.90				$^{13}\text{C}$	4.32	1.63,1.63	1.24,1.19	2.86
	55.00	33.50	24.46	28.57	40.90						55.52	32.35	23.96	41.06

<sup>a</sup> $^1\text{H}$  and  $^{13}\text{C}$  chemical shifts are reported relative to 3-(trimethylsilyl)propionic- $d_4$  acid. <sup>b</sup>There are two sets of chemical shifts for Pro-2 corresponding to the N-terminal Ala-1 and the des-Ala-1 forms of IL-1 $\beta$  (Driscoll et al., 1990a). P2\* refers to the des-Ala-1 form.

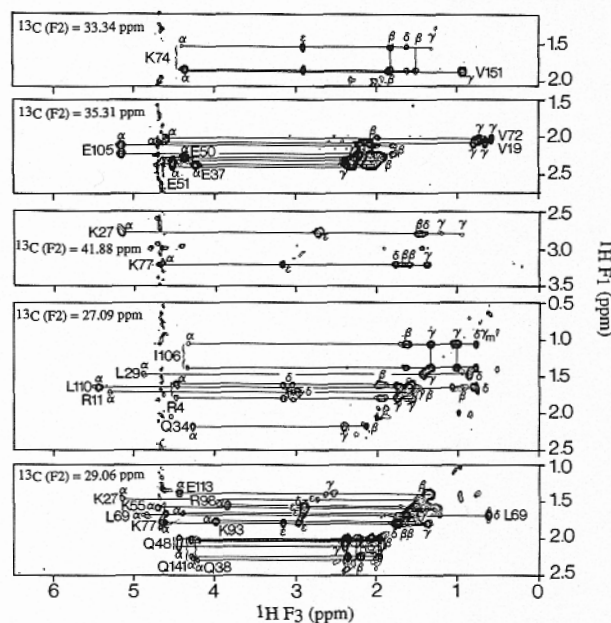


FIGURE 10: Selected slices at different  $^{13}\text{C}(F_2)$  chemical shifts of the HCCH-TOCSY spectrum of uniformly  $^{15}\text{N}/^{13}\text{C}$ -labeled IL-1 $\beta$  illustrating connectivities originating from the  $\text{C}^\beta\text{H}$  and  $\text{C}^\gamma\text{H}$  protons of Glu and Gln, the  $\text{C}^\gamma\text{H}$  protons of Arg, the  $\text{C}^\beta\text{H}$ ,  $\text{C}^\delta\text{H}$ , and  $\text{C}^\epsilon\text{H}$  protons of Lys, and the  $\text{C}^\gamma\text{H}$  protons of Ile.

to be significantly enhanced for exterior side chains due to an increase in mobility as one proceeds outward along the Arg side chains. In contrast, intensities of cross-peaks to the  $\text{C}^\delta\text{H}$  of proline would be expected to be reduced on account of the motional restrictions imposed by the five-membered proline ring. This is indeed what is observed: the  $\text{C}^\alpha\text{H}-\text{C}^\delta\text{H}$  connectivity for Arg-98 at  $\delta^{13}\text{C} = 57.99$  ppm is clearly stronger than the  $\text{C}^\alpha\text{H}-\text{C}^\beta\text{H}$  and  $\text{C}^\alpha\text{H}-\text{C}^\gamma\text{H}$  cross-peaks, while for Pro-23 and Pro-87 at  $\delta^{13}\text{C} = 63.57$  ppm the  $\text{C}^\alpha\text{H}-\text{C}^\delta\text{H}$  cross-peak intensities are comparable or weaker than the other two. In the same slice as these two Pro residues, one can also see the complete Arg-4 spin system originating from the  $\text{C}^\delta\text{H}$  protons at  $\delta^{13}\text{C} = 42.87$  ppm, with cross-peak intensities following a characteristic order,  $\text{C}^\delta\text{H}-\text{C}^\alpha\text{H} > \text{C}^\delta\text{H}-\text{C}^\gamma\text{H} \sim \text{C}^\delta\text{H}-\text{C}^\beta\text{H}$ , in agreement with theoretical expectations (Figure

5). Finally, complete Lys spin systems originating from the  $\text{C}^\alpha\text{H}$  proton are also clearly seen in Figure 7 (cf. Lys-94 at  $\delta^{13}\text{C} = 53.38$  ppm, Lys-55 and Lys-74 at  $\delta^{13}\text{C} = 55.03$  ppm, Lys-93 at  $\delta^{13}\text{C} = 58.65$  ppm, and Lys-97 at  $\delta^{13}\text{C} = 59.30$  ppm). Note how easily the  $\text{C}^\alpha\text{H}-\text{C}^\beta\text{H}$ ,  $\text{C}^\alpha\text{H}-\text{C}^\delta\text{H}$ , and  $\text{C}^\alpha\text{H}-\text{C}^\gamma\text{H}$  relayed cross-peaks are observed. In this respect it is worth pointing out that, even in small proteins of 40–60 residues, it is rare to be able to delineate complete lysine, arginine, and proline spin systems originating from the  $\text{C}^\alpha\text{H}$  proton by use of conventional 2D  $^1\text{H}-^1\text{H}$  HOHAHA spectroscopy, as even there the ratios of line widths to  $^1\text{H}-^1\text{H}$  three-bond couplings are unfavorable for efficient magnetization transfer.

As already noted, the interpretation of the HCCH-COSY and HCCH-TOCSY spectra is rendered essentially foolproof through extensive cross-checks afforded by examining the connectivities between  $^1\text{H}$  resonances observed at all the corresponding  $^{13}\text{C}$  shifts of the directly bonded  $^{13}\text{C}$  nuclei. This procedure is illustrated in Figure 8, which shows a series of relayed connectivities originating from the  $\text{C}^\delta\text{H}$  (10.66 and 12.96 ppm),  $\text{C}^\gamma\text{H}$  (15.26 and 16.57 ppm), and  $\text{C}^\beta\text{H}$  (39.91 and 40.90 ppm) of four Ile residues (Ile-56, -104, -106, and -122). In virtually every case, the complete Ile spin system is identified from each of these starting points, thus confirming the  $^1\text{H}$  assignments derived from the connectivities originating at the  $\text{C}^\alpha\text{H}$  protons and providing the  $^{13}\text{C}$  chemical shifts of the other carbon atoms. It is worth pointing out here too that while in small proteins the  $\text{C}^\beta\text{H}$  and  $\text{C}^\gamma\text{H}$  resonances are relatively easily identified in 2D  $^1\text{H}-^1\text{H}$  COSY and HOHAHA spectra, connectivities to the  $\text{C}^\gamma\text{H}$  and  $\text{C}^\delta\text{H}$  resonances are often difficult to observe owing to large unresolved multiplet widths of the  $\text{C}^\gamma\text{H}$  resonances. In the HCCH-TOCSY spectrum, however, assignment of these protons is a simple matter.

A similar series of connectivities for Leu residues originating from the  $\text{C}^\delta\text{H}$  methyl (slices at  $\delta^{13}\text{C} = 22.16$ , 23.80, and 24.79 ppm) and  $\text{C}^\gamma\text{H}$  protons (slice at  $\delta^{13}\text{C} = 25.45$  ppm) is shown in Figure 9. It may be noticed that the  $\text{C}^\delta\text{H}$  to  $\text{C}^\beta\text{H}$  relayed peaks and the direct  $\text{C}^\gamma\text{H}-\text{C}^\beta\text{H}$  peaks tend to be rather weak and sometimes even absent, in complete agreement with the simulations shown in Figure 5, so that the HCCH-COSY experiment (or alternatively a short 8–10-ms mixing time



HCCH-TOCSY experiment) still has a valuable role to play in confirming the identification of the Leu C $\beta$ H resonance positions.

Finally, a number of connectivities originating successively from the C $\beta$ H ( $\delta^{13}\text{C}$  = 33.34 ppm), C $\alpha$ H ( $\delta^{13}\text{C}$  = 29.06 ppm), and C $\gamma$ H ( $\delta^{13}\text{C}$  = 41.88 ppm) protons of Lys and from the C $\gamma$ H proton of Arg ( $\delta^{13}\text{C}$  = 27.09 and 29.06 ppm) and the C $\gamma$ H protons of Glu ( $\delta^{13}\text{C}$  = 35.31 ppm) are shown in Figure 10. As in the previous examples, complete or virtually complete spin systems can be delineated starting from the different-side chain protons, thereby providing multiple checks on the assignments.

It is clear from the above examples that the combined use of the 3D HCCH-COSY and HCCH-TOCSY spectra provides a very powerful means for the unambiguous assignment of both the  $^1\text{H}$  and  $^{13}\text{C}$  chemical shifts of the side-chain protons and carbons. Using these spectra, we were able to obtain complete  $^1\text{H}$  and  $^{13}\text{C}$  assignments for the aliphatic side chains of all residues with the exception of only four lysine residues (Lys-16, Lys-63, Lys-92, and Lys-103) for which partial assignments were obtained. A complete list of assignments is provided in Table I.

The absence of complete connectivities for four of the lysine residues may be due to multiple side-chain conformations or to line broadening reducing the efficiency of magnetization transfer. Line broadening may arise from conformational restriction of lysine side-chain positions (cf. the lysine residue in the bottom panel of Figure 4 with nonequivalent C $\beta$ H as well as C $\gamma$ H resonances, which additionally attenuates cross-peak intensities) or from an intermediate exchange process. The latter may well be the case for Lys-63 as the  $^{15}\text{N}$  line width of the backbone amide group of this residue is extensively broadened by a chemical exchange contribution to  $T_2$  (Clare et al., 1990b).

For completeness we have also included in Table I the  $^1\text{H}$  chemical shifts of the aromatic protons together with the  $^{13}\text{C}$  shifts of their directly bonded carbons. The  $^1\text{H}$  aromatic spin systems had previously been delineated by 2D  $^1\text{H}$ - $^1\text{H}$  correlation methods (P.COSY and HOHAHA) and sequence specifically assigned from the observation of characteristic intrarésidue NOEs from the aromatic ring protons to the NH, C $\alpha$ H, and C $\beta$ H protons (Driscoll et al., 1990a). With these assignments in hand, the assignments of the  $^{13}\text{C}$  shifts of the aromatic carbons attached to protons involved in intraaromatic NOEs could be obtained in a trivial manner from a 3D  $^1\text{H}$ - $^{13}\text{C}$  NOESY-HMQC spectrum (data not shown). The latter experiment was necessary as the present HCCH-COSY and HCCH-TOCSY experiments in which the carbon carrier was placed at 43 ppm could not be used to detect correlations involving aromatic carbons owing to the fact that the aromatic carbons resonate at the downfield end of the  $^{13}\text{C}$  spectrum at  $\sim 130$  ppm. As a result, adequate refocusing of  $^{13}\text{C}$  chemical shifts of aromatic carbons by the  $^{13}\text{C}$  180° pulses before and after the  $^{13}\text{C}$  90° COSY mixing pulse was not obtained in the case of the HCCH-COSY experiment (cf. pulse sequence under Experimental Procedures), and the rf field strength used in the HCCH-TOCSY experiment was too weak to permit efficient isotropic mixing of the aromatic carbons.

The aliphatic  $^{13}\text{C}$  side-chain assignments provided in Table I provide a reasonably large and comprehensive database of  $^{13}\text{C}$  chemical shifts in proteins and expand on the previously available data on isolated amino acids and small peptides (Levy & Nelson, 1972; Howarth & Lilley, 1978) and the partial assignments recently obtained for staphylococcal nuclease by 2D heteronuclear methods (Wang et al., 1990). The distri-

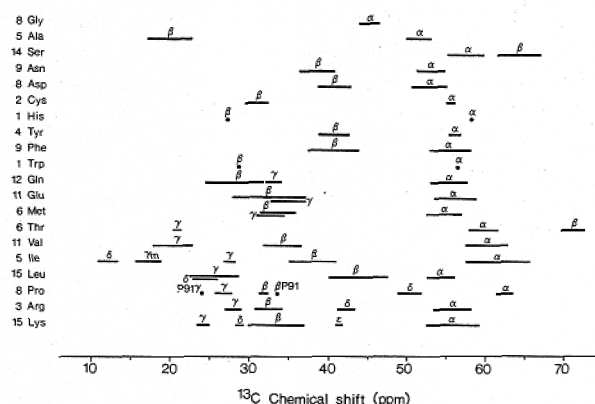


FIGURE 11: Summary of the  $^{13}\text{C}$  chemical shift ranges observed for the different carbon atoms and different residue types in IL-1 $\beta$ . The number of residues present in IL-1 $\beta$  for each amino acid type is indicated at the left-hand side of the figure. In the case of the Pro residues, there are two sets of chemical shifts for Pro-2 corresponding to the N-terminal Ala-1 and des-Ala-1 forms of IL-1 $\beta$  (Driscoll et al., 1990a). Complete  $^{13}\text{C}$  assignments were obtained for 11 out of the 15 lysine residues; no  $^{13}\text{C}$  shifts were obtained for Lys-63, only  $^{13}\text{C}\alpha$  shifts were obtained for Lys-92 and Lys-103, and only  $^{13}\text{C}\alpha$ ,  $^{13}\text{C}\beta$ , and  $^{13}\text{C}\epsilon$  shifts were obtained for Lys-16.

bution of  $^{13}\text{C}$  aliphatic chemical shifts in IL-1 $\beta$  for the different amino acid residues is summarized graphically in Figure 11. The chemical shifts are comparable to those established for isolated amino acids and small peptides, but not surprisingly, the ranges for each carbon type are significantly larger. The information provided in Figure 11 may therefore serve as a useful aid in future assignment work. In general, chemical shifts are usually difficult to interpret. Nevertheless, empirical correlations may be of some use if treated with appropriate caution. For example, it has been noted in peptides that the C $\alpha$  chemical shift of the residue preceding a proline is  $\sim 2$  ppm upfield from its characteristic value (Torchia et al., 1975). There are eight prolines in IL-1 $\beta$  at positions 2, 23, 57, 78, 87, 91, 118, and 131. In five cases, the C $\alpha$  chemical shift of the preceding residue (Ile-56, Lys-77, Asp-86, Phe-117, and Met-130) has the most upfield chemical shift in its residue class, while in another two (Ala-1 and Gly-22) the C $\alpha$  chemical shift is clearly at the upfield end of its range. In the case of Tyr-90, which precedes Pro-91, there is no clear correlation. This, however, may have to do with the fact that the Tyr-90-Pro-91 peptide bond is in the *cis* conformation as evidenced from the observation of characteristic C $\alpha$ H-C $\alpha$ H sequential NOEs between Tyr-90 and Pro-91 in the  $^1\text{H}$ - $^{13}\text{C}$  HMQC-NOESY spectrum, in agreement with the crystallographic findings for the X-ray structure of IL-1 $\beta$  (Finzel et al., 1989). The other proline residues, on the other hand, are all in the *trans* conformation as judged from both the NMR and X-ray data. In this regard, it is worth noting that, just as in the case of small peptides (Howarth & Lilley, 1978), the C $\beta$  and C $\gamma$  chemical shifts of the *cis*-proline at position 91 are about 1 ppm downfield and 2 ppm upfield, respectively, of the corresponding shifts for the *trans*-proline residues.

#### CONCLUDING REMARKS

The determination of protein structures in solution by NMR depends on the identification of a very large number of short ( $< 5$  Å) approximate interproton distance restraints derived from NOE measurements [see Wüthrich (1986, 1989), Clare and Gronenborn (1987, 1989), and Gronenborn and Clare (1990) for reviews]. A key requirement in this procedure involves not simply the assignment of the backbone  $^1\text{H}$  resonances (NH and C $\alpha$ H) but equally importantly that of the side-chain resonances. Even for proteins of 50–100 residues,

it is often difficult to obtain complete side-chain assignments for long side chain amino acids such as Lys and Arg. This is not simply due to spectral overlap but to the relatively large  $^1\text{H}$  line widths which render magnetization transfer through homonuclear  $^1\text{H}$ - $^1\text{H}$  couplings inefficient. Further, even when these assignments are made, extensive spectral overlap often precludes their use in the identification of NOESY cross-peaks. The data presented in this paper clearly show that 3D heteronuclear methods based on magnetization transfer via large one-bond  $^1\text{H}$ - $^{15}\text{N}$ ,  $^{15}\text{N}$ - $^{13}\text{C}$ ,  $^1\text{H}$ - $^{13}\text{C}$ , and  $^{13}\text{C}$ - $^{13}\text{C}$  couplings are highly efficient and sensitive even for molecules the size of IL- $1\beta$ , which comprises 153 residues. In this respect we note that all the 3D spectra were recorded on a single 1.7 mM sample of uniformly  $^{15}\text{N}/^{13}\text{C}$ -labeled protein and the limiting factor in the duration of the experiments was not the signal-to-noise but the minimum number of phase cycles required to reduce spectral artifacts to an acceptable level. Indeed, all the 3D heteronuclear correlation spectra for the present paper were recorded in just 1 week of measuring time. Further, with the addition of a third dimension afforded by the  $^{13}\text{C}$  nucleus, problems associated with extensive resonance overlap which are invariably present in the 2D  $^1\text{H}$ - $^1\text{H}$  and  $^1\text{H}$ - $^{13}\text{C}$  spectra of proteins the size of IL- $1\beta$  are completely overcome.

To our knowledge, the  $^1\text{H}$ ,  $^{15}\text{N}$ , and  $^{13}\text{C}$  assignments of IL- $1\beta$  provided by this paper and the preceding one (Driscoll et al., 1990a) represent the first example of a protein larger than 15 kDa for which essentially complete assignments of both backbone and side-chain resonances have been obtained by heteronuclear 3D NMR methods. Given that the attainable precision of a protein structure determination depends crucially on maximizing the number of approximate interproton distance restraints (Driscoll et al., 1989a,b; Kraulis et al., 1989), virtually complete assignments represent the necessary prerequisite in that endeavor. With the assignments for IL- $1\beta$  firmly in hand, the way is now open for the full determination of a high-resolution 3D structure of IL- $1\beta$  in solution by 3D  $^1\text{H}$ - $^{15}\text{N}$  and  $^1\text{H}$ - $^{13}\text{C}$  HMQC-NOESY spectroscopy, as well as 4D  $^{13}\text{C}/^{15}\text{N}$  and  $^{13}\text{C}/^{13}\text{C}$  edited NOESY spectroscopy (Kay et al., 1990c), at a precision comparable to the best that is now attainable for smaller proteins (Kraulis et al., 1989; Qian et al., 1989; Clore et al., 1990a).

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